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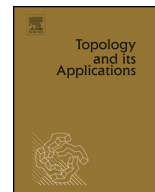
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A geometrical take on invariants of low-dimensional manifolds found by integration



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ABSTRACT

An elementary geometrical proof of the fact that the Euler characteristic is the only topological invariant of a surface that can be found by integration (using Gauss–Bonnet) is given. A similar method is also applied to three-dimensional manifolds.

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1. Introduction

The Gauss–Bonnet theorem relates the integral of some intrinsic quantity whose origins lie in the field of differential geometry, namely the Gaussian curvature, to some topological invariant, the Euler characteristic. For higher-dimensional manifolds the Gauss–Bonnet theorem can be generalized, using the theory of characteristic classes. For a very elegant exposition we refer to Milnor and Stasheff [12] or alternatively Spivak [13]. Abrahamov [1] proved that the invariants thus produced are unique, up to some equivalence. See Gilkey [7] for a modern (and more extensive) treatment. Below we provide a proof of a similar statement for two- and three-dimensional manifolds, based solely on geometrical arguments, in contrast to the more algebraic approach taken in the literature.

The formulation of the main result will be along the lines of the following question proposed by I.M. Singer: ‘Suppose that f is a scalar valued invariant of the metric such that $t(M) = \int f \, \text{dvol}$ is independent of the metric. Then is there some universal constant c so that $t(M) = c\chi(M)$?’ This question has reportedly [7] been answered in the affirmative by E. Miller.

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The proof as discussed by Gilkey is somewhat algebraic in nature and focuses on the invariance of $f = F(g, \partial g, \dots)$ under coordinate transformations. Fortunately the functions which are well behaved can be easily found and be listed, using a theorem by Weyl on the invariants of the orthogonal group. The invariant functions thus found are linear combinations of contractions of Riemann tensors and their derivatives. The functions can be further distinguished based on their behavior under rescaling of the metric. If the product of the function and the volume form is invariant under this rescaling it is a candidate for a topological invariant.¹ For example in two dimensions the Gaussian curvature is the only such function (up to some remainder whose integral is zero). It can be shown that all such functions yield topological invariants.

Our geometrical proof relies heavily on the classification of two-dimensional closed surfaces and on Heegaard splitting. A discussion of the classification can be found in [9] or [10], for the latter we refer to [5] or [14]. We complete our discussion by some remarks on generalizations and the implications of this result in the study of the large scale structure of the universe, which caused our interest.

2. Surfaces

Theorem 1. *Let f be a function on two-dimensional real Riemannian compact manifolds, which is completely determined by the metric, in the sense that f locally can be written as $f(x) = F(g(x), \partial g(x), \dots)$ where g denotes the metric, independent of the topology of the base manifold. Suppose the integration*

$$I_f(M) \equiv \int_M f \, \text{dvol},$$

where dvol indicates the volume form, of f over an orientable² manifold M yields a topological invariant $t_f(M)$ for all surfaces. We write $t_f(M)$ to emphasize the dependence on f . Then there exists a real number c_f , depending only on f , such that $t_f(M) = c_f \chi(M)$, where χ denotes the Euler characteristic.

Proof. First we note that the space of Riemannian metrics on a manifold is connected. This is obvious because if g and \tilde{g} are metrics then so is $\lambda g + (1 - \lambda)\tilde{g}$ for all $\lambda \in [0, 1]$. This means that we can assume without loss of generality that M is isometrically embedded in \mathbb{R}^3 , because we can choose \tilde{g} to be the standard metric of M . Now let f be a function as described in the theorem, such that

$$\int_M f \, \text{dvol} = t$$

is a topological invariant. Suppose that for the two-sphere \mathbb{S}^2 we have

$$\int_{\mathbb{S}^2} f \, \text{dvol} = 2c,$$

where c is some constant. From this we can conclude that for the sphere $t = c\chi(M)$.

We can now deform the two-sphere as follows. A small region is pushed outwards and bent – in a sufficiently smooth manner – such that this region contains three equally spaced parallel cylinder pieces all of the same radius. We can now cut in the cylindrical part along the plane orthogonal to the cylinder and reassemble the parts so that we recover a topological sphere but also get a torus. The integral is not altered because integrals are additive. The procedure is illustrated in Fig. 1. Because the integral is clearly additive for unions this implies that

¹ A line of reasoning one also encounters in the work by Abrahamov.

² Clearly the integral over a non-orientable manifold does not make sense.

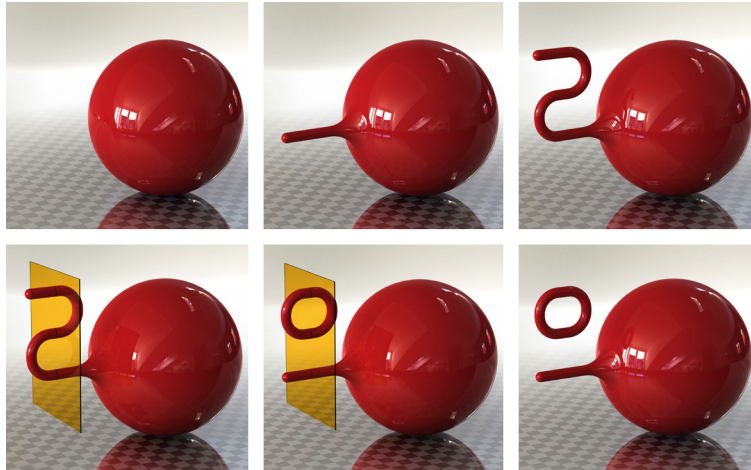


Fig. 1. From left to right, top to bottom we have depicted: the sphere, the deformation (in two steps), the deformed surface with cutting lines indicated by the yellow glass plane and the reassembled surfaces (with and without cutting lines). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

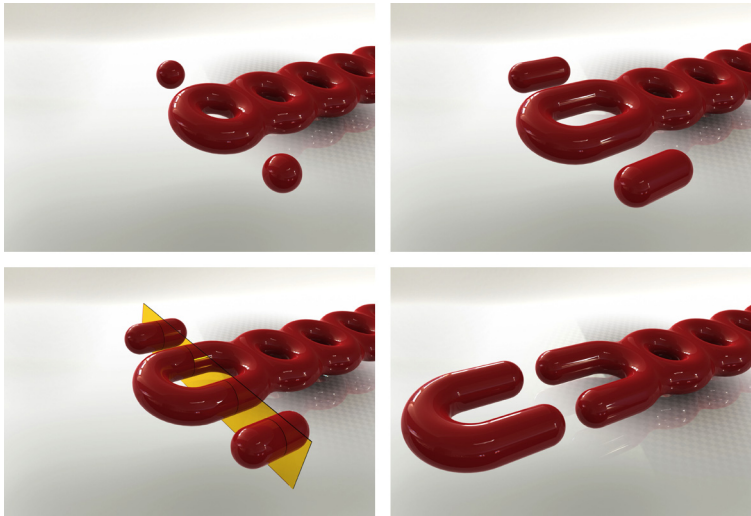


Fig. 2. From left to right, top to bottom we have depicted: An n -torus and two spheres, the same surfaces deformed, the deformed surfaces with the lines along which we cut indicated by yellow glass plane and the reassembled surfaces. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

$$\int_{C_1} f \, d\text{vol} = 0,$$

where C_1 is a surface of genus 1. Generally we shall denote a surface of genus g by C_g .

The rest of the proof is inductive in nature. We begin with a topological genus- g torus and two spheres. We deform these surfaces so that the spheres contain a piece of a cylinder, both of the same radius, and the n -torus such that it contains two pieces of the cylinder, again of the same radius, so that if these pieces are deleted one of the remaining surfaces is itself a topological cylinder. We again cut the cylindrical pieces in half and reassemble the parts so that we have a genus- $g - 1$ torus and a sphere, as sketched in Fig. 2.

We can now conclude that

$$\int_{C_g} f \, d\text{vol} + 2 \int_{\mathbb{S}^2} f \, d\text{vol} = \int_{C_{g-1}} f \, d\text{vol} + \int_{\mathbb{S}^2} f \, d\text{vol}$$

and thus by induction that

$$\int_{C_g} f \, \text{dvol} = c(2 - 2g) = c\chi(C_g).$$

By the classification of all 2-manifolds we have proven the theorem for all two-dimensional real manifolds embedded in \mathbb{R}^3 . \square

Remark 2. In [Theorem 1](#) we assumed that f gives us a topological invariant for all surfaces, in fact the conclusion can be drawn for a given manifold M , if $\int f \, \text{dvol}$ is an invariant for S^2 , $S^1 \times S^1$ and M .

The proof of this statement differs from that above in that instead of the induction step illustrated in [Fig. 2](#), we consider a genus- g surface and $2g$ balls and perform the cut and paste operation for each hole simultaneously.

3. Three dimensions

We will now focus on the three-dimensional case. The intuition for the following proof is much strengthened by the remark that a Morse function h on some manifold M can always be interpreted as height function. This can be easily seen as follows: Let M be isometrically embedded in \mathbb{R}^n , possibly using the Nash embedding theorem. Then we can add the value of the Morse function as another coordinate to a point $p \in M \subset \mathbb{R}^n$, so that the manifold M is embedded in \mathbb{R}^{n+1} and the last coordinate is the height.

Theorem 3. *Let f be a function on three-dimensional real Riemannian compact manifolds, which is completely determined by the metric, in the sense that f locally can be written as $f(x) = F(g(x), \partial g(x), \dots)$ where g denotes the metric, independent of the topology of the base manifold. If the integration*

$$I_f(M) \equiv \int_M f \, \text{dvol}$$

of f over a manifold M yields a topological invariant $t_f(M)$, for all 3-manifolds, then we have $t(M) = 0$.

Proof. The first step in our proof will consist of showing that if $M = C_g \times S^1$ we have that

$$\int_M f \, \text{dvol} = 0.$$

To show this we shall consider a manifold N , that admits a Heegaard splitting of genus g . This means that the manifold N can be represented as the attachment of two three-dimensional manifolds, which are both homeomorphic to a three-dimensional ball with g handles, with respect to a diffeomorphism of their boundaries. We further have that there exists a Morse function h on N with one minimum and one maximum and all critical points of index 1, 2 correspond to the critical values c_1 and c_2 respectively with $c_1 < c_2$, see [\[5\]](#). This has been schematically represented in the leftmost picture in [Fig. 3](#).³

³ Note that conversely a Heegaard splitting also gives a Morse function in a natural manner. Namely we start with Morse functions on both g -handled balls, by simply taking a Morse function on the standard g -handled ball and pulling back via the diffeomorphisms to the g -handled balls in question. Now Theorem 1.4 and Lemma 3.7 of [\[11\]](#) give a differentiable structure on the union with a smooth structure compatible with the given differentiable structure on the different parts, moreover such that the Morse functions on both parts piece together to a smooth function.

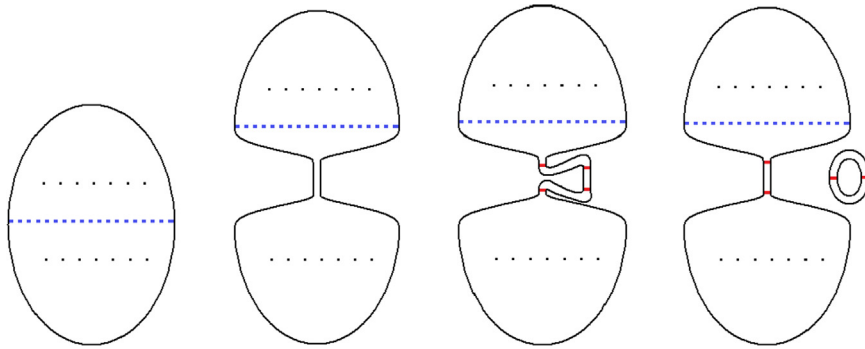


Fig. 3. From left to right we have sketched: A manifold admitting a Heegaard splitting; the critical points of the Morse function are indicated as dots and the attachment by a blue dashed line, the same manifold with a small part of it brought to a standard $C_g \times [-\delta, \delta]$ metric, the deformed surface with cutting lines (red) indicated and the reassembled surfaces. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

We now define for every surface C_g of genus g , some metric induced by an embedding in \mathbb{R}^3 , exhibiting \mathbb{Z}_2 symmetry. We shall refer to this Riemannian manifold as the standard surface of genus g . In the following we view N as embedded in \mathbb{R}^k . Let f be as in Theorem 1 such that

$$\int_N f \, d\text{vol},$$

is a topological invariant t . For some sufficiently small $[a_1, b_1] \subset \mathbb{R}$, with $c_1 < a_1 < \alpha_1 < \beta_1 < b_1 < c_2$, we smoothly and isotopically deform $h^{-1}([a_1, b_1]) \cap M \sim C_g \times [a_1, b_1]$, so that $h^{-1}([\alpha_1, \beta_1]) \cap M$ becomes isometric to the standard $C_g \times [\alpha_1, \beta_1] \subset \mathbb{R}^4 \subset \mathbb{R}^k$ given by the standard C_g and the ordinary Cartesian product. We shall now deform this part of the manifold so that it consists of a straight piece and two pieces which are straight at the beginning and the end but are bent in the middle so that if we cut along the boundaries of the pieces and reassemble we recover the original manifold and $C_g \times \mathbb{S}^1$. The procedure is sketched in Fig. 3. From this we conclude that

$$\int_N f \, d\text{vol} = \int_N f \, d\text{vol} + \int_{C_g \times \mathbb{S}^1} f \, d\text{vol},$$

where we again used local isotopy and the additivity of integration. Therefore,

$$\int_{C_g \times \mathbb{S}^1} f \, d\text{vol} = 0.$$

The next part of the proof relies on the fact that the sphere (\mathbb{S}^3) allows a Heegaard splitting of every genus g , see [5]. Let M be a manifold which allows a Heegaard splitting of genus g . We now deform two pieces of the manifold into parts isometric to $C_g \times [\alpha_1, \beta_1]$ and $C_g \times [\alpha_2, \beta_2]$, with $\alpha_1 < \beta_1 < \alpha_2 < \beta_2$, so that for all $p_1 \in (\alpha_1, \beta_1)$ and $p_2 \in (\alpha_2, \beta_2)$ both $h^{-1}((-\infty, p_1)) \cap M$ and $h^{-1}(p_2, \infty) \cap M$ are topological spheres with g handles whose boundary is isometric to the standard genus- g surface, as discussed above. We can now smoothly deform $h^{-1}((p_1, q_1)) \cap M$ and $h^{-1}((q_2, p_2)) \cap M$, with $p_1 < q_1 < \beta_1$ and $\alpha_2 < q_2 < p_2$ (see Fig. 5), such that if we cut along the p_i and q_i lines and reassemble (possibly using \mathbb{Z}_2 symmetry) we recover two topological manifolds, with given topology. One of the manifolds we thus construct is a manifold admitting a Heegaard splitting of genus g . The attachment diffeomorphism, of the latter, on the boundary of the sphere with g handles is the identity. This manifold shall be denoted by $M_g^{S(3D)}$. The other manifold is a topological $C_g \times \mathbb{S}^1$ -manifold. The entire procedure is sketched in Fig. 4.

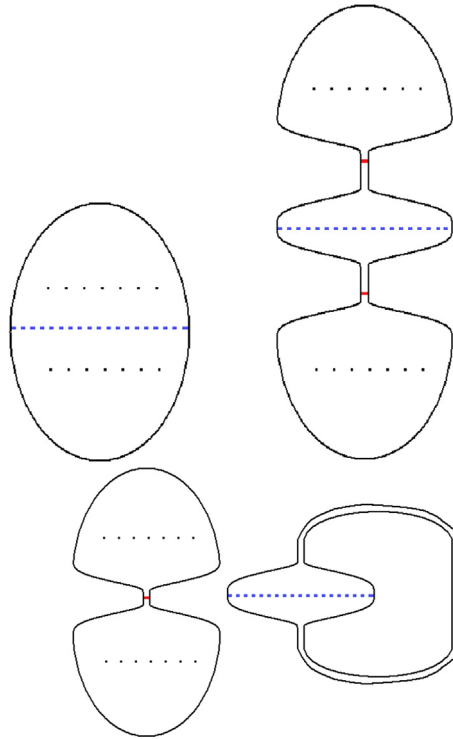


Fig. 4. A Heegaard splitting, then the same manifold with two small parts brought to a standard metric both on another side of the ‘attachment line’, cutting lines (red) are also indicated, and finally the reassembled surface (two connected components). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

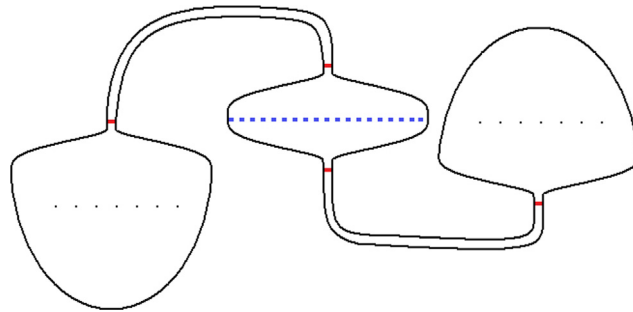


Fig. 5. A sketch of the deformed manifold with the cutting lines (red) and the ‘attachment line’ (blue). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

This means that by deforming, cutting and pasting a manifold M , which allows a Heegaard splitting of genus g , we find the following equalities

$$\int_M f \, \text{dvol} = \int_{M_g^{S(3D)}} f \, \text{dvol} + \int_{C_g \times \mathbb{S}^1} f \, \text{dvol} = \int_{M_g^{S(3D)}} f \, \text{dvol} + 0,$$

where f is as defined in the theorem. If we now use that the sphere (\mathbb{S}^3) allows a Heegaard splitting of every genus g we find that

$$\int_M f \, \text{dvol} = \int_{M_g^{S(3D)}} f \, \text{dvol} = \int_{\mathbb{S}^3} f \, \text{dvol}. \quad (1)$$

Following this observation, we are able to use the result of the first part of the proof,

$$\int_{C_g \times \mathbb{S}^1} f \, \text{dvol} = 0.$$

This immediately translates into

$$\int_{\mathbb{S}^2 \times \mathbb{S}^1} f \, \text{dvol} = 0.$$

We notice that both \mathbb{S}^3 and $\mathbb{S}^2 \times \mathbb{S}^1$ allow a Heegaard splitting of genus 1, so that

$$\int_{\mathbb{S}^3} f \, \text{dvol} = \int_{M_1^{S(3D)}} f \, \text{dvol} = \int_{\mathbb{S}^2 \times \mathbb{S}^1} f \, \text{dvol} = 0. \quad (2)$$

Combining Eqs. (1) and (2) yields

$$\int_M f \, \text{dvol} = 0,$$

for any manifold M and $f = f(g, \partial g, \dots)$ a function determined by the metric and all its derivatives. \square

Remark 4. In [Theorem 3](#) we assumed that f gives us a topological invariant for all 3-manifolds, in fact the conclusion can be drawn for a given manifold M , if $\int f \, \text{dvol}$ is an invariant S^3 , $S^1 \times C_g$, $M_1^{S(3D)}$ and M , where C_g is a surface of genus g and $M_1^{S(3D)}$ as defined above.

This is clear from inspection of the proof of [Theorem 3](#).

4. Discussion

One can wonder about generalizations of the methods stated above to manifolds of general dimension. Some of these generalizations are immediately obvious, for example the procedure sketched in [Fig. 3](#) can be used in any dimension so see that for f and t as in the theorem

$$\int_{M^{d-1} \times \mathbb{S}^1} f \, \text{dvol} = t$$

implies that $t = 0$, where M^{d-1} is any manifold of dimension $d - 1$ occurring as level set. However a full classification of all integrals yielding a topological invariant does not seem feasible because there is no easy classification of manifolds of dimension $d - 1$ for $d > 3$ (and none for $d > 4$), occurring as the level sets of a Morse function on a manifold of dimension d .

Recently a lot of effort has been put into the study of topological properties of Gaussian random fields [\[2,6,15\]](#), using among others Euler integration [\[4\]](#). This is especially important in the context of cosmology, because these random fields are believed to describe the density fluctuations in the early universe [\[3\]](#). A nice closed formula for the expectation value of the Euler characteristic of the level sets of these fields has been found [\[3,8\]](#), because it could be determined using the Gauss–Bonnet theorem. The result above shows that no expressions can be found for all other interesting topological invariants, such as Betti numbers, associated to the field, using a similar straightforward integration technique.

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